

# Lecture 15 The Lagrangian Floer TFT, II

Recall:  $(M, \omega)$  symplectic  $\mathcal{J} = \{\omega\text{-compatible a.s.s.}\}$

$S = \hat{S} \setminus (\Sigma^- \cup \Sigma^+)$  pointed boundary Riemann surface with strip like ends  $\varepsilon_s: \mathbb{R}^{\pm} \times [0, 1] \rightarrow S$  for  $s \in \Sigma^{\pm}$ .

A set of Lagrangian labels  $\{L_c\}_{c \in \pi_0(\partial S)}$   
 Floer data  $(H_s \in C^{\infty}([0, 1], \mathcal{H}), J_s \in C^{\infty}([0, 1], \mathcal{J}))$   
 for each pair  $(L_{s,0}, L_{s,1})$ .

compatible perturbation data  $(K \in \Omega^1(S, \mathcal{H}), J \in C^{\infty}(S, \mathcal{J}))$   
 $Y = \text{Ham.v.f.}(K) \in \Omega^1(S, C^{\infty}(TM))$

$$\mathcal{E}(L_{s,0}, L_{s,1}) = \{y: [0, 1] \rightarrow M \mid y(i) \in L_{s,i} \ (i=0,1), \dot{y} = X_{H_s}(t, y(t))\}$$

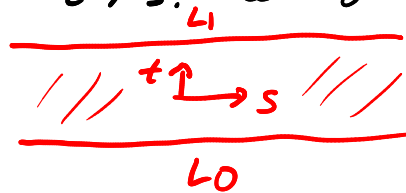
Moduli space  $\mathcal{M}(\{y_s\}_{s \in \Sigma}) = \text{solutions } u: S \rightarrow M \text{ of}$

$$\begin{cases} (du - Y)^{0,1} = 0 \\ u(c) \in L_c \text{ for } c \in \pi_0(\partial S) \\ \lim_{s \rightarrow \pm\infty} u(\varepsilon_s(s, t)) = y_s(t) \text{ for } s \in \Sigma^{\pm} \end{cases}$$

These moduli spaces are used to define the TFT operation.

First, consider the case  $S = \mathbb{Z} = \mathbb{R} \times [0, 1]$ . Let  $L_0$  and  $L_1$  be the Lagrangian labels.

Choose Floer datum  $(H, J)$  for  $(L_0, L_1)$ . Then take



an S-TRANSLATION INVARIANT Perturbation datum  $(K, J')$

$$K = H(t) dt \quad J'(s, t) = J(t).$$

The equation  $(du - Y)^{0,1} = 0$  is more explicitly written

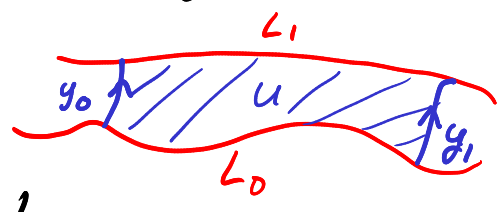
$$\begin{aligned} \partial_s u + J(t, u) (\partial_t u - X(t, u)) &= 0 \\ u(s, 0) \in L_0 \quad u(s, 1) \in L_1 \end{aligned}$$

The form of the equation is unchanged by translation.

in the  $s$ -direction. Write  $\mathcal{M}_Z(y_0, y_1)$  for the set of solutions satisfying the asymptotic conditions

$$\lim_{s \rightarrow +\infty} u(s, t) = y_1(t) \quad \lim_{s \rightarrow -\infty} u(s, t) = y_0(t),$$

where  $y_0, y_1 \in \mathcal{C}(L_0, L_1)$



Then  $\mathcal{M}_Z(y_0, y_1)$  has an  $\mathbb{R}$ -action, and

we denote by  $\mathcal{M}_Z^*(y_0, y_1) = \mathcal{M}_Z(y_0, y_1) / \mathbb{R}$

Let  $K$  be a field of  $\text{char}(K) = 2$ .

Define the Floer cochain group  $CF^{pr}(L_0, L_1)$  to be a  $K$ -vector space with a basis vector for each  $y \in \mathcal{C}(L_0, L_1)$

$$CF^{pr}(L_0, L_1) := \bigoplus_{y \in \mathcal{C}(L_0, L_1)} K \cdot y$$

The differential is defined in terms of the basis as

$$\partial(y_1) = \sum_{y_0} \# \mathcal{M}_Z^*(y_0, y_1) y_0$$

where  $\#$  means "count the isolated points modulo 2".

More generally, for any pointed boundary R.S.  $S = \hat{S} \setminus (\Sigma^- \cup \Sigma^+)$

$$\mathcal{M}_S(\{y_s\}_{s \in \Sigma}) = \left\{ \begin{array}{l} (du - \gamma)^{0,1} = 0 \\ u(c) \subset L_c \\ \lim_{s \rightarrow \pm\infty} u(E_s(s, t)) = y_s(t) \end{array} \right\}$$

Define  $C\Phi_S : \bigotimes_{s^+ \in \Sigma^+} CF^{pr}(L_{s^+, 0}, L_{s^+, 1}) \rightarrow \bigotimes_{s^- \in \Sigma^-} CF^{pr}(L_{s^-, 0}, L_{s^-, 1})$

in terms of basis:  $C\Phi_S(\bigotimes_{s^+} y_{s^+}) = \sum_{\{y_{s^-}\}} \# \mathcal{M}_S(\{y_{s^-}, y_{s^+}\}) (\bigotimes_{s^-} y_{s^-})$

- Now there are some things we would like to be able to claim
- Desiderata: (I) on  $CF^{\text{pr}}(L_0, L_1)$ ,  $\partial\partial = 0$ , so  $\partial$  really is a differential, and  $CF^{\text{pr}}(L_0, L_1)$  is a complex and we can take its cohomology  $HF^{\text{pr}}(L_0, L_1)$
- (II)  $C\Phi_S$  is a chain map, so it induces a map  $\Phi_S: \bigotimes_{S^+} HF^{\text{pr}}(L_{S^+,0}, L_{S^+,1}) \rightarrow \bigotimes_{S^-} HF^{\text{pr}}(L_{S^-,0}, L_{S^-,1})$
- (III)  $HF^{\text{pr}}(L_0, L_1)$  is invariant under change of Floer data
- (IV)  $\Phi_S$  is invariant under change of perturbative data, and change of conformal structure on  $S$ .
- (V)  $\Phi_S$  satisfies the gluing axiom of a TFT.

These points are all somewhat intertwined, and in order to address them we first need to understand better how moduli spaces like  $\mathcal{M}_S(\{y_S\})$  are actually constructed. Furthermore, simple examples show these desiderata cannot always be obtained without some restrictions on  $(M, \omega)$  and/or  $L_C$ , we will clarify this as we go.

Outline of construction of  $\mathcal{M}_S(\{y_S\})$  as a manifold (not just a set)

Let  $\mathcal{B}_S^\infty \subset \text{Map}(S, M)$  be the space of smooth maps  $u: S \rightarrow M$  such that  $u(C) \subset L_C$  and which converge to some collection  $\{y_S\}$  on the strip-like ends.

This is an  $\infty$ -dimensional "manifold" and its tangent space at  $u$   $C^\infty((S, \partial S), (u^*TM, u^*TL_C)) \subset C^\infty(S, u^*TM)$  consisting of sections  $\xi$  of  $u^*TM$  over  $S$  such that, along the component  $C \subset \partial S$ ,  $\xi \in u^*TL_C$ .

over  $\mathcal{B}_S^\infty$  there is a natural  $\infty$ -rank vector bundle  $E_S^\infty$  whose fiber at  $u \in \mathcal{B}_S^\infty$  is  $C^\infty(S, \Omega_S^{0,1} \otimes u^*TM)$

The mapping  $\bar{\partial}_S: u \mapsto (du - Y)^{0,1}$  can then be regarded as a section:

$$\begin{array}{c} E_S^\infty \\ \downarrow \pi \\ \mathcal{B}_S^\infty \end{array} \Bigg) \bar{\partial}_S \quad \text{and } \mathcal{M}(\{y_S\}) \text{ is (a component of) vanishing set of this section.}$$

To give  $\mathcal{M}(\{y_S\})$  a manifold structure, we need to be able to formulate a condition of transversality in this  $\infty$ -dimensional setting, and require  $\bar{\partial}_S$  to be transverse to the zero section.