

Physics $N=2$ SCFT \rightsquigarrow A or B twist.
(2,2)

$$Q = \begin{cases} Q_A \\ Q_B \end{cases} \Rightarrow \mathcal{H}_{dual} = \mathcal{H}_{A,B}^{(def)} = \langle \phi_i \rangle_{i=1}^5$$

\uparrow
deformation subspace

$$Q^2 = 0$$

parametric marginal deformations
(deformation is still SCFT)

$$A(H^{3,1}(X)) = \mathcal{H}_A^{def}$$

$$B(H^{3,1}(X)) = \mathcal{H}_B^{def}$$

for $(\mathbb{C}P^3)$

Marginal Def from ϕ_i ?

$$S \rightarrow S + \int_{\mathbb{C}\Sigma} d^2z \phi_i^{(2)}$$

$$\phi_i^{(2)} := \left\{ G^{\bullet}, [G^{\bullet}, \phi_i] \right\}$$

$\uparrow \quad \uparrow$
super-currents

$N=2$ SCFT \rightsquigarrow Vertex Alg (OPE's)

$$\oint G^-(z) \rightarrow \mathcal{Q}^- \quad \widehat{U(1)}_L \times \widehat{U(1)}_R$$

~~$[Q, A(z)]$~~

$[Q, A(z_2, \bar{z}_2)] = \text{Res}_{z_1 \rightarrow z_2} j(z_1) A(z_2, \bar{z}_2)$

holomorphic change \uparrow \uparrow any op \uparrow \uparrow ope \uparrow ope
 $Q = \oint \frac{dz}{2\pi i} j(z)$

$= \oint_{C_{z_2}} j(z_1) A(z_2, \bar{z}_2)$

To keep action real, do:

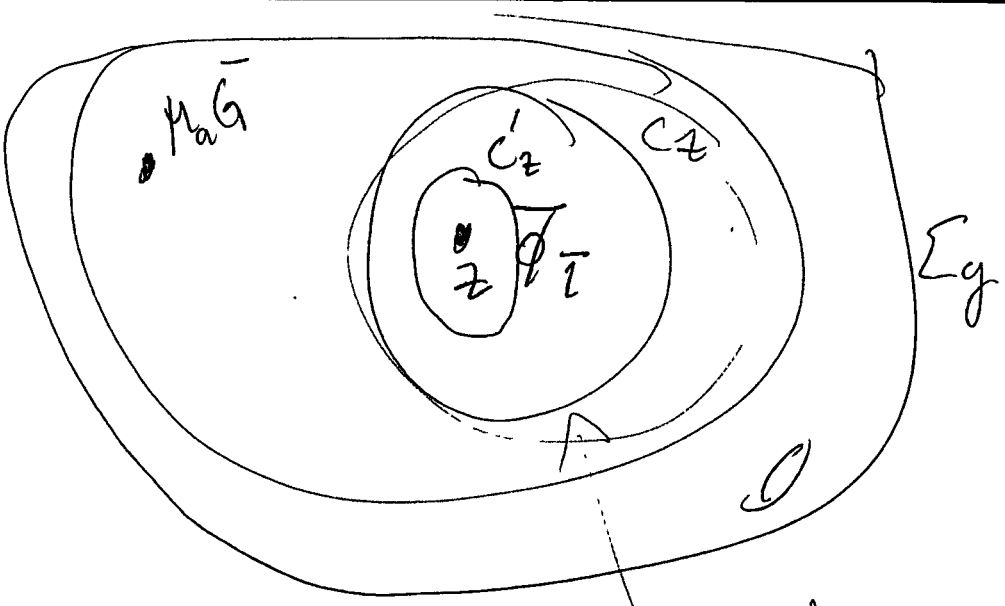
$S \rightarrow S + t^i \int_{\Sigma} d^2z \phi_i^{(2)} + \bar{t}^i \int_{\Sigma} d^2z \bar{\phi}_i^{(2)}$

$F_g = \int_{\mathcal{M}_g} [dm] \left\langle \prod_{a=1} \int \mu_a G^- \int \bar{\mu}_a G^- \right\rangle_{\Sigma_g}$

where $\int \mu_a G^- := \int_{\Sigma_g} \mu^{(a)}_{\bar{z}} G_{z\bar{z}}^- dz_1 d\bar{z}$

$\bar{\partial}_{\bar{z}} F_g = \int_{\mathcal{M}_g} [dm] \int_{\Sigma} d^2z \left\langle \oint_{C_2} G^+ \oint_{C'_2} G^+ \bar{\phi}_i(z) \prod (\dots) \right\rangle$

$\parallel \frac{\partial}{\partial t^i}$



can be seen as encircling z ,
 or as encircling the complement
 with opposite orientation.

If $C_z \rightsquigarrow$ surrounds $z \Leftrightarrow \oint_{C_z} G^+ \cdot \theta$

Do residue calculus.

$$G^+(z)G^-(w) \text{ contains } \frac{2T(w)}{(z-w)} + \text{more singular}$$

$T(w)$ Stress energy tensor. (classically $T_{\alpha\beta} \sim \frac{\delta S}{\delta g^{\alpha\beta}}$)
 (modes = Virasoro)
 metric $\sim \Sigma_g$

Two copies of $N=2$ SCA. (G^+, G^-, T, \dots)
 they commute with $(\bar{G}^+, \bar{G}^-, \bar{T}, \dots)$
 each other.

~~$$\oint_{C_z} G^+(w) dw \int dz \mu_{\bar{z}}^{(b)} G_{zz}(z)$$~~

Commutator / get $2 \int M_b^T$
use OPE

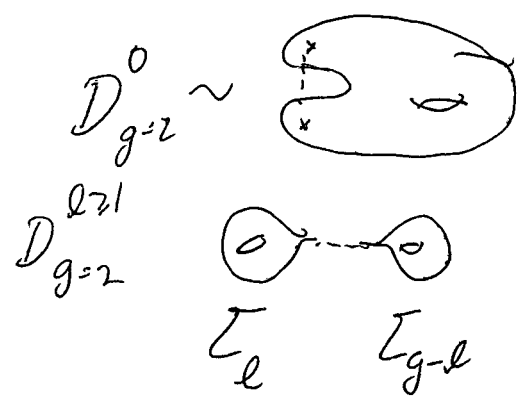
$$\int_{M_g} [dm] \sum_{b=1}^{3g-3} \left\langle \int_{\bar{\mathcal{D}}_1} \int 2\mu_b^T \int 2\bar{\mu}_b^T \prod_{a \neq b} \int M_a^- G^- \int M_a^- G^- \right\rangle_{\Sigma_g}$$

Next track: $\int 2\mu_b^T \int 2\bar{\mu}_b^T \Leftrightarrow 4 \frac{\partial^2}{\partial m_b \partial \bar{m}_b}$
 m_b coords on M_g

~~$\int_{M_g} [dm]$~~

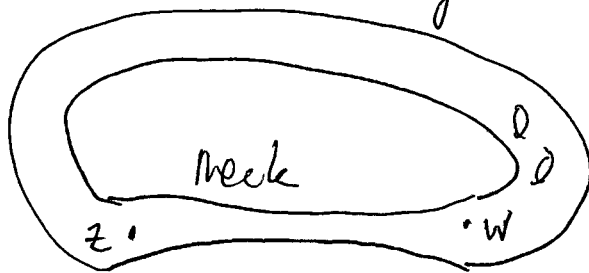
$$\int_{M_g} [dm] \sum_{b, \bar{b}} \frac{4 \partial^2}{\partial m_b \partial \bar{m}_b} \left\langle \int_{\bar{\mathcal{D}}_1} \prod M_a^- G^- \int M_a^- G^- \right\rangle$$

$$\partial M_g = \bigsqcup_{l=0}^{[g/2]} D_g^{(l)}$$



Contrib of D_g^0

Close to D_g^0



moduli $m = (m', z, w, \tau)$

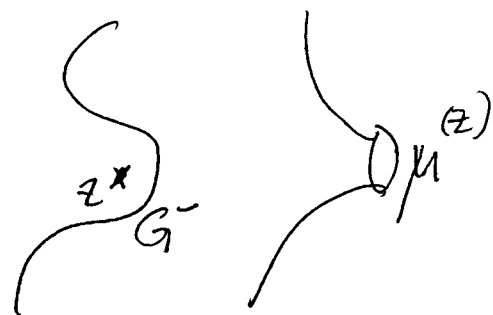
$\begin{matrix} | & \backslash & \backslash & \backslash \\ \text{coordinates} & \text{points} & & \text{(length, torsion)} \\ \text{on } \mathcal{M}_{g-1} & & & \end{matrix}$

$$\int_{\mathcal{M}_g} [dm] \quad [dm', dw, dz, d\tau]$$

$\frac{\partial^2}{\partial m_b \partial \bar{m}_b}$
contains
 $\frac{\partial^2}{\partial \tau \partial \bar{\tau}}$
 \rightsquigarrow
 $\frac{\partial}{\partial \text{Im}(\tau)}$

$$\int_{\Sigma_g} \mu^{(z)} G^- \rightsquigarrow \int_{C_z} G^-$$

↑ associated to marked part z

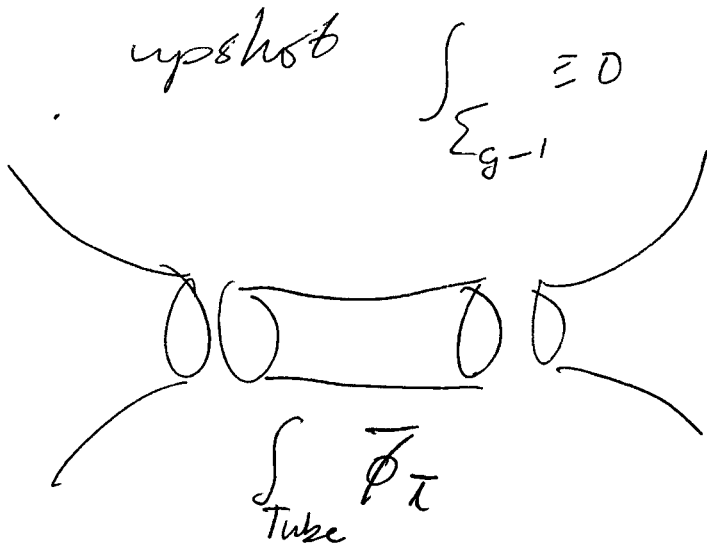


$$\int_{D_g^0} [dm] [dw, dz] \frac{\partial}{\partial \text{Im} \tau} \left(\int_{\Sigma_g} \bar{\phi}_i \int_{C_z} G^- \int_{C'_z} G^- \int_{C_w} G^- \int_{C'_w} G^- \right)$$

$$\prod_{a=1}^{3g-6} \int_{\Sigma_{g-1}} M'_a G^- \int_{\Sigma_{g-1}} \bar{M}_a G^-$$

This is essentially a residue.

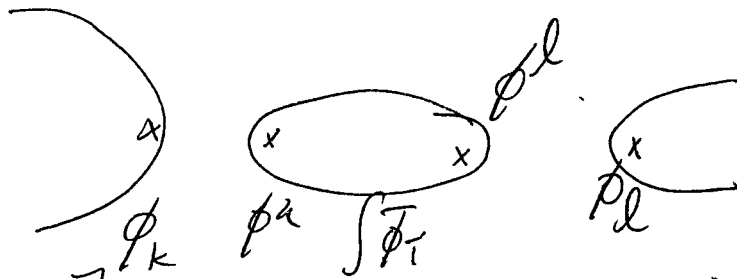
$$\int_{\Sigma_g} \bar{\phi}_i = \int_{\Sigma_{g-1}} \bar{\phi}_i + \int_{\text{tube}} \bar{\phi}_i$$



metric

$$\langle \phi_i, \phi_j \rangle = \langle \phi_i, \phi_j \rangle_{\mathcal{H}}$$

$$= \eta_{ij}$$



$$\langle \phi_k, \phi^J \rangle = \delta_k^J$$

only Σ_k are summed over

$$\rightarrow \phi_i(z) \eta^{ii'} \langle i' | \int \bar{\phi}_i | k' \rangle \eta_{kk'}$$

G_{ref}

~~independent~~

$$\langle j | \bar{\phi}_i | k \rangle$$

Kähler pot + metric

[7]

$$= \bar{C}_{\bar{j}\bar{k}} e^{2k} G^{\bar{j}\bar{j}'} G^{\bar{k}\bar{k}'} \eta_{j\bar{j}} \eta_{k\bar{k}}$$

independent

of where $\bar{\phi}_i$ is, so

$$\{G^-[G, \bar{\phi}]\}$$

$$\bar{C}_{\bar{j}\bar{k}} e^{2k} G^{\bar{j}\bar{j}'} G^{\bar{k}\bar{k}'} \int_{\mathcal{M}_{g-1}} [dm'] \left(\int_{\Sigma_{g-1}} \phi_j^{(2)} \int_{\Sigma_{g-1}} \phi_k^{(2)} \right)$$

$$\frac{3g-6}{\prod_{a=1} \int \mathcal{M}'_a G^- \int \mathcal{M}'_0 G^- \Bigg\rangle_{\Sigma_{g-1}}$$

~~$\langle \int \phi_j^{(2)} \int \phi_k^{(2)} \prod_{a=1} \int \mathcal{M}'_a G^- \int \mathcal{M}'_0 G^- \rangle_{\Sigma_{g-1}}$~~

$$\left\langle \int \phi_j^{(2)} \int \phi_k^{(2)} \prod_{a=1} \int \mathcal{M}'_a G^- \int \mathcal{M}'_0 G^- \right\rangle_{\Sigma_{g-1}}$$

||

$$D_j D_k \mathcal{F}_{g-1}$$

cov. of \mathcal{D} .