

# Si Li BCOV theory and Quantized B-model

1

Outline: ① BCOV theory and period map.

② Quantization

③ Elliptic curve example

$X$ : CY geometry (compact CY/LG)

$\Omega_X$ : CY volume form.

$$PV(X) = \Omega^{0,0}(\Lambda^0 T_X^{1,0}) = \bigoplus_{i,j} PV^{i,j} \quad \text{"} \Omega^{0,d}(\Lambda^i T_X) \text{"}$$

$$PV^{i,j} \xleftrightarrow{+\Omega_X} \Omega^{d-i,j}$$

$$\bar{\partial}, \partial \leftrightarrow \bar{\partial}, \partial$$

$$\partial : PV^{i,j} \rightarrow PV^{i-1,j} \quad \text{determined w/r/t } \Omega_X$$

$$- \{ \alpha, \beta \} = \partial(\alpha\beta) - (\partial\alpha)\beta - (-1)^\alpha \alpha\partial\beta$$

$$- Q = \bar{\partial} + t\partial$$

Period map: Assume  $X \subset \mathbb{C}P^3$

13

Consider  $\text{Def}(X, \Omega_X) \rightarrow H^3(X_0, \mathbb{C})$

$[X, \Omega_X] \rightarrow [\Omega_X]$

•  $H^3(X, \mathbb{C})$  is a symplectic space  $\int_X \alpha \wedge \beta$

•  $\text{Im } P$  is a Lagrangian (Linear)

• Let  $L \subset H^3$  be Lagrangian s.t.

$$H^3 = \mathbb{P}^2 \oplus L$$

$$\cong (H^{3,0} \oplus H^{2,1})$$

$$\text{dim } H^3 \cong T^*\mathbb{P}^2$$

$$\text{Im}(P_X) = \text{Graph}(dF_0)$$

$$F_0 \in \mathcal{O}(H^{3,0} \oplus H^{2,1})$$

isomorphic to  $\mathbb{C}W^2$   
 $g=0$

Diff  $L$   
 leads to def theory.

Remark:  $\text{Def}(X, \Omega_X) \quad \mu \in \mathbb{P}V^{1,1} \quad \bar{\partial}\mu + \frac{1}{2}\{\mu, \mu\} = 0$

$$d(\mathcal{O}e^n + \Omega_X) = 0 \quad \text{CY volume form}$$

$$\Leftrightarrow Q(\mu + t\rho) + \frac{1}{2}\{\mu + t\rho, \mu + t\rho\} = 0$$

$Q = \bar{\partial} + D \quad t$  formal variable

<sup>u</sup> cyclic cohomology describes  $\text{def}(X, \Omega_X)$

3

Def<sup>n</sup>: The fields of BCOV theory

$$(S_+ = PV(X)[[t]], \quad Q = \bar{\partial} + t\partial)$$

$$\text{let } S_- = t^{-1}PV(X)[[t^{-1}]]$$

$$S = PV(X)((t))$$

$$\Gamma(S, Q) \longrightarrow \Omega^\bullet(X)((t))$$

$$t^k \alpha \longrightarrow t^{k+i-1} \alpha \in \Omega_X, \alpha \in PV^{i,*}$$

$$Q \longrightarrow d$$

$$\Gamma(S_+) = \prod_{p \in \mathbb{Z}} t^{d-p-1} \Omega^{\geq p,*}$$

$$S_+ \subset S \iff \text{Hodge filtration.}$$

Symplectic Pairing on  $S$

$$\omega(f(t)\alpha, g(t)\beta) = \left( \text{Res}_{t=0} f(t)g(t)dt \right) \text{Tr}(\alpha\beta)$$

$$\text{Tr} : PV \rightarrow \mathbb{C} \quad \mu \mapsto \int_X (\mu \lrcorner \Omega_X) \lrcorner \Omega_X$$

$\Rightarrow (S, Q, \omega)$  is symplectic space.

4

Remark  $X = CY_3$  sub symplectic space

$$t PV^{0,0} \oplus t PV^{1,1} \oplus t^{-1} PV^{2,2} \oplus t^{-2} PV^{3,3}$$



$$\Omega^{3,0} \oplus \Omega^{2,1} \oplus \Omega^{1,2} \oplus \Omega^{0,3}$$

Descendants  
Hodge wt

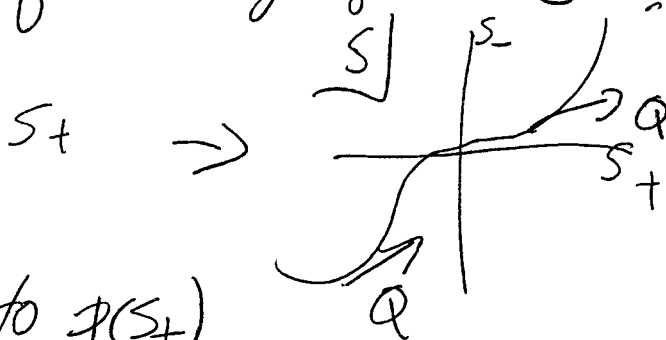
Lagrangian cone (after Givental)

$$\mathcal{P}: S_+ \rightarrow S$$

$$\mu \rightarrow t e^{\mu/t} - t$$

as a formal scheme ( $\mu$  nilpotent)

①  $\mathcal{P}(S_+)$  is a formal Lagrangian  $\subset S$



②  $Q$  is tangent to  $\mathcal{P}(S_+)$

③  $\pi_+ : S \rightarrow S_+ \quad \mathcal{P}(S_+) \xrightarrow{\pi} S_+$   
formally diff.

④  $\mathcal{P}_*(Q, \xi, \zeta) = Q$   
push-forward nonlinear vector field, get linear one.

means

$$Q(t e^{N/t}) = \left( Q_\mu + \frac{1}{2} \{ \mu, \mu \} \right) e^{N/t} \quad (\text{BTT}) \quad [5]$$

Defn: BCOV interaction  $I_X^{\text{BCOV}}$  is a local functional on  $S_+ = \text{PV}(X)[[t]]$

$$\mathcal{P}(S_+) = \text{Graph} (dI_X^{\text{BCOV}})$$

(viewed for  $S = T^*S_+$ )

Lemma [Costello-Li]  $I_X(\mu) = \text{Tr} \langle e^{\mu} \rangle_0$

$$\langle t^{k_1} \mu_1, \dots, t^{k_n} \mu_n \rangle_0$$

$$= \left( \int_{\overline{M}_{0,n}} \psi_1^{k_1} \dots \psi_n^{k_n} \right) \mu_1 \dots \mu_n$$

const map  $(\mathbb{P}^1, X) \simeq \overline{M}_{0,n} \times X$

Degenerate BV theory

$$S = \begin{array}{ccc} S_+ & \xleftarrow{Q} & \mathbb{Q} \\ \oplus & \uparrow & \\ S_- & \xrightarrow{Q} & \mathbb{Q} \end{array}$$

Failure of  $Q$  to preserve splitting at cochain level

$\Downarrow$   
degenerate Poisson.

Let  $\omega^{-1}$  be the Poisson kernel

(6)

$$\omega^{-1} = \sum_k (-t)^k \otimes t^{-k-1} \delta_0$$

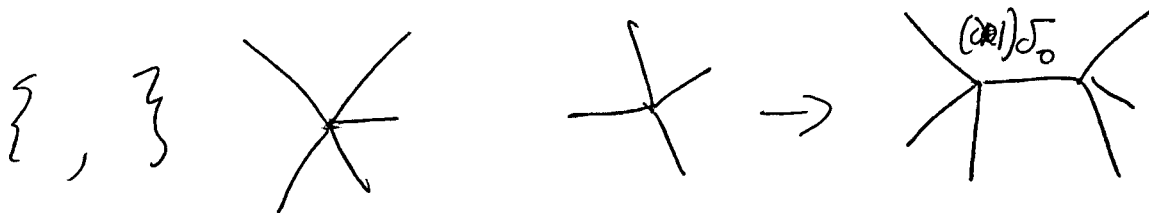
$$K_0 = \pi_+ \otimes \pi_+ ((Q \otimes 1) \omega^{-1}) \text{ Poisson kernel.}$$

$$= (\partial \otimes 1) \delta_0 \quad Q = \bar{\partial} + \partial$$

degenerate because splitting not preserved.

BV bracket

$$\mathcal{O}_{loc}(S_+) \otimes \mathcal{O}_{loc}(S_+) \rightarrow \mathcal{O}_{loc}(S_+)$$



Lemma:  $I_X^{BCOV}$  satisfies CME

$$Q I_X^{BCOV} + \frac{1}{2} \{ I_X^{BCOV}, I_X^{BCOV} \} = 0$$

Geometrically

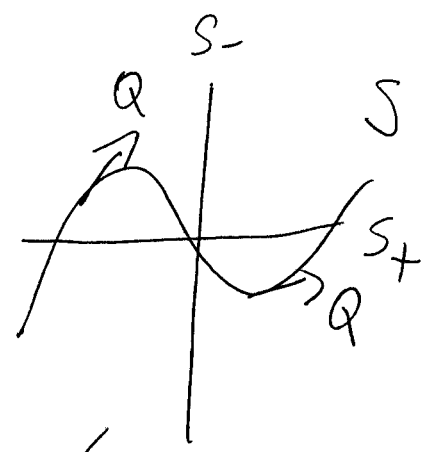
$$\pi_{+*}(Q) = Q + \{ I_X^{BCOV}, - \} \text{ defines } L_\infty \text{ str in } S_+$$

Lox-quasi iso

$$S_+ \xrightarrow{\pi_+ \circ P} S_+$$

$$Q, \xi, \beta \rightarrow Q + \{I_x, -\}$$

$$\frac{Q, \xi, \beta}{S_+} \xrightarrow{P}$$



$$\mu \rightarrow [te^{\mu t} - t]_+$$

↑  
No  $\beta$  variable

↑  
is variational

$$\frac{Q + \{I_x, -\}}{S_+} \xleftarrow{\pi_+}$$

Quantization Toy Model  $(V, \omega, Q)$

$Q \rightsquigarrow V_+ \subset V$  linear lagrangian subcomplex  
 $\omega^{-1}$ , Poisson kernel.

Defn:  $\mathcal{W}(V) = \Pi (V^*)^{\otimes n} [h] / \sim$

$$[a, b] = a \star b - b \star a \sim h \omega^{-1}(a, b)$$

$a, b \in V^*$

$$\text{Ann}(V_+) \subset V^* \text{ the annihilator of } V_+$$

$$= \{ \varphi : V \rightarrow \mathbb{C} \mid \varphi(V_+) = 0 \}$$

$$\implies \text{Fock}(V) := \mathcal{W}(V) / \mathcal{W}(V) \text{Ann}(V_+)$$

$$\psi$$

$$|0\rangle \leftrightarrow 1$$

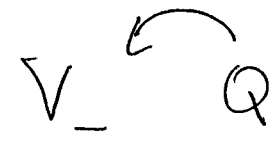
$Q$  preserves  $V$  and  $\omega^{-1}$  and  $V_+$

$\Rightarrow (W(V), Q), (\text{Fock}(V), Q)$

Quantization: a state ~~is~~

$|F\rangle \in \text{Fock}(V)$  st.  $Q|F\rangle = 0$

Let  $V = V_+ \oplus V_-$  isotropic splitting.



$\Rightarrow V^* = V_+^* \oplus \text{Ann}(V_+)$

$Q \curvearrowright \mathcal{O}(V_+) \simeq \widehat{\text{Sym}}(V_+^*) \rightarrow W(V) \otimes Q$

doesn't preserve  $Q$

$\downarrow$   
 $\text{Fock}(Q) \otimes Q$

Defn:  $I \in \mathcal{O}(V_+)[\hbar]$  satisfies QME

$(Q + \hbar \partial_k) e^{I/\hbar} = 0$

$e^{I/\hbar} \leftrightarrow |I\rangle = e^{I/\hbar} |0\rangle$

QME:  $QI + \hbar \partial_k I + \frac{1}{2} \{I, I\} = 0 \Leftrightarrow Q|I\rangle = 0$



# Quantum BCOV theory

19

$$(S = S_+ \oplus S_-, Q, \omega)$$

$$K_0 = (\partial \otimes 1) \delta$$

$\mathcal{O}(S_+)$  = distributions

$\partial_{K_0} \rightsquigarrow \mathcal{O}(S_+)$  ill-defined.

homotopic RG flow

$\rightsquigarrow$

$|I_r\rangle$  vary homotopically

$\rightsquigarrow$

convergence, doesn't merge etc.

To RG flow,  
 $|I\rangle$  changes  
by  $Q$ -exact.

$\leftarrow$

Can calculate this  
at any scale!

( $r \rightarrow \infty$  is easier, maybe)

$$|I\rangle \in \text{Fock}(H^*(S_+, Q))$$

Get numbers: Choose a splitting  $H^*(S, Q) = H^*(S_+, Q) \oplus L$

$$\Rightarrow \text{Fock}(H^*(S_+, Q)) \stackrel{L}{=} \mathcal{O}(H^*(S_+, Q))$$

$$\nearrow H^*(S_+ = PV[\mathbb{C}P^1], Q) \stackrel{L}{\cong} H^*(\mathbb{R}P^1, \mathbb{C})$$

Splitting on

conformal is ok, but need  $L$  for convergence

Ell curve:  $E = \mathbb{C} / (2\pi i \mathbb{Z})$

(10)

$$S_t = \underbrace{\Omega_{E}^{0,k}}_{b_k t^k} [t] \oplus \underbrace{\Omega_{E}^{0,k}}_{\eta_k t^k} (T_E [t]) [t]$$

Stationary sector:  $b_{>0} = 0$      $\eta = \text{const}$

$$I_X = \sum_{k \geq 0} \int_E \left( \eta_k \frac{b_0^{k+2}}{(k+2)!} \right) \quad (\text{CME}_E) \quad \text{sur.}$$

CME  $\Leftrightarrow$  classical integrable hierarchy

QME  $\Leftrightarrow$  Quantum integrable hierarchy

$$\Rightarrow I_X^q = \sum_{k \geq 0} \int \eta_k \frac{W^{(k+2)}(b_0)}{k+2}$$

$$W^{(k+2)}(b_0) = \sum \frac{k_i!}{\prod k_i!} \left( \prod_i \frac{1}{i!} (\sqrt{h} \partial_z)^{i-1} b_0 \right)^{k_i}$$

scale zero (UV finite)

$\Rightarrow$  generating function

$$T_r \text{ Heisenberg vertex } q \quad e^{\sum \eta_k \int \frac{W^{(k+2)}}{k+2}} \quad q = \exp(2\pi i t)$$

minor to  $\langle t^{k_1} \omega_1, t^{k_2} \omega_2, \dots, t^{k_n} \omega_n \rangle$

(11)

$\text{Tr } e^{-t\Delta}$
$S^1 \rightarrow X$

observe:  $h_4(x, y) \sim \int D\gamma e^{-|\dot{\gamma}|^2}$

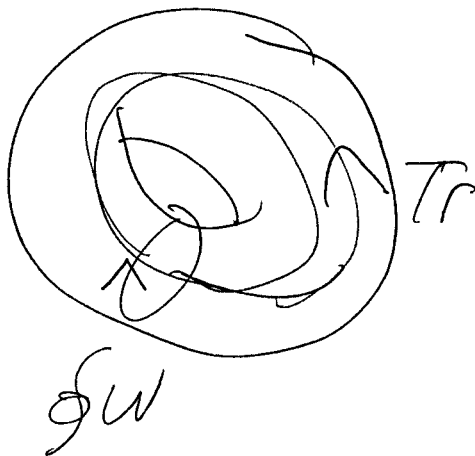
$$\gamma: [0, 1] \rightarrow X$$

$$\gamma(0) = x$$

$$\gamma(1) = y$$

$$\text{Trace} = \int_{\text{loop space}} e^{-|\dot{\gamma}|^2}$$

2d



$$9 \text{ } L_0^{-\frac{1}{24}}$$

analyse of  $e^{-t\Delta}$