

## Riemannian Geometry II: Levi-Civita connection.

Now we can reap the fruits of our labor in setting up the theory of connections. For any Riemannian manifold  $(M, g)$ , there is a canonical connection in the tangent bundle  $TM$ , which is determined uniquely by the metric  $g$ .

Let  $(M, g)$  be a Riemannian manifold. We consider connections in the tangent bundle  $TM$ , thought of as a covariant derivative  $\nabla$ . Recall that  $\mathcal{X}(M) = \Gamma(M, TM)$  is the set of vector fields on  $M$ . Thus a connection in  $TM$  is a map

$$\begin{aligned} \nabla : \mathcal{X}(M) \times \mathcal{X}(M) &\longrightarrow \mathcal{X}(M) \\ (X, Y) &\longmapsto \nabla_X Y \end{aligned}$$

Satisfying the properties we set down before.

- (i)  $\nabla_X(aY + bZ) = a\nabla_X Y + b\nabla_X Z$  ( $a, b \in \mathbb{R}$ )
- (ii)  $\nabla_X(fY) = (X.f)Y + f\nabla_X Y$  ( $f \in C^\infty(M)$ )
- (iii)  $\nabla_{(fX + gZ)} Y = f\nabla_X Y + g\nabla_Z Y$  ( $f, g \in C^\infty(M)$ )

The neat thing about connections in  $TM$  (as opposed to connections in a general vector bundle  $E$ ) is that the space of sections of  $TM$  is the space of vector fields, so  $\nabla$  is letting us differentiate a vector field  $Y$  with respect to another vector field  $X$ , and the result is a third vector field  $\nabla_X Y$ .

How could  $\nabla$  be related to the metric  $g$ ?

Let's go back to  $\mathbb{R}^n$ . Let  $\vec{v}$  and  $\vec{w}$  be vector fields on  $\mathbb{R}^n$ , thought of as functions  $\vec{v}, \vec{w} : \mathbb{R}^n \rightarrow \mathbb{R}^n$ . It is certainly true that

$$\underbrace{\frac{\partial}{\partial x^i} (\vec{v} \cdot \vec{w})}_{\text{This is the derivative of a function.}} = \underbrace{\frac{\partial \vec{v}}{\partial x^i}}_{\text{These are derivatives of vector fields}} \cdot \vec{w} + \vec{v} \cdot \underbrace{\frac{\partial \vec{w}}{\partial x^i}}_{\text{These are derivatives of vector fields}}$$

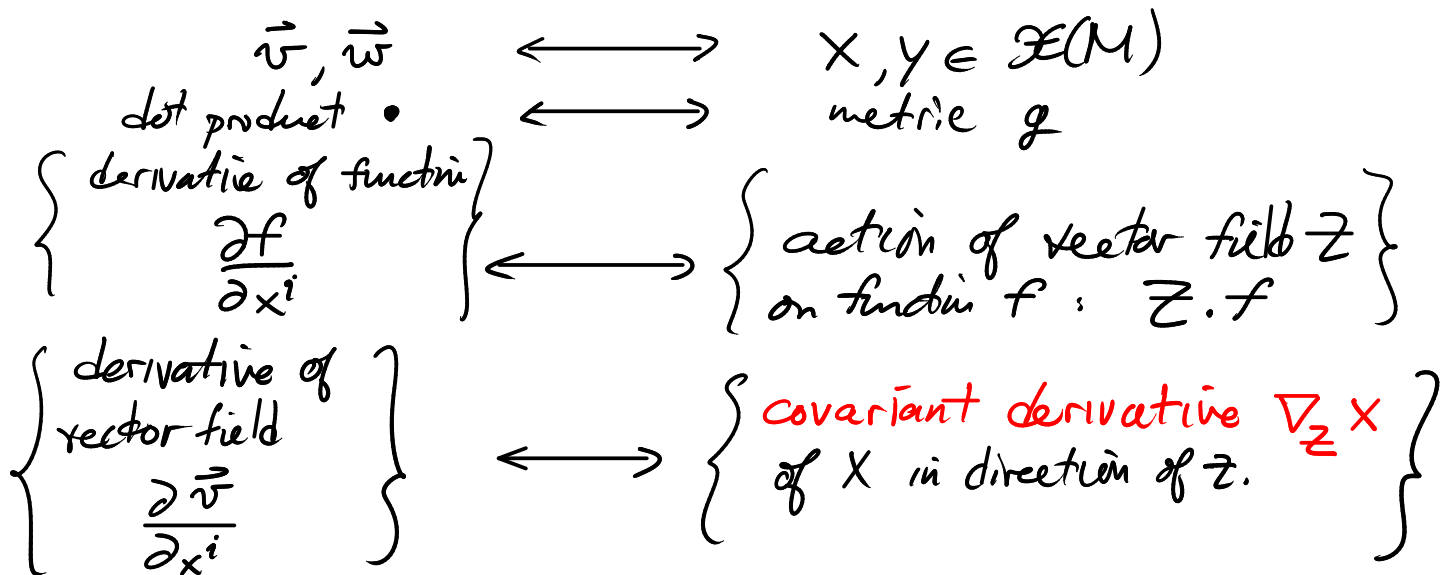
This is the derivative of a function.

These are derivatives of vector fields

We would like to generalize this to a Riemannian manifold  $(M, g)$

Standard  $\mathbb{R}^n$

$(M, g)$



Thus, the identity we want is written

$$\otimes Z \cdot (g(X, Y)) = g(\nabla_Z X, Y) + g(X, \nabla_Z Y)$$

Not every connection  $\nabla$  satisfies this! If  $g$  is given, this is a condition on  $\nabla$ .

Def A connection  $\nabla$  in TM is compatible with  $g$  if the identity  $\otimes$  is satisfied by all vector fields  $X, Y, Z \in \mathcal{X}(M)$ .

Remark: This definition makes sense for a connection in any vector bundle equipped with a Euclidean metric.

For connections in the tangent bundle TM, there is another special property we can consider: Recall the definition of Lie bracket, in  $\mathbb{R}^n$  or in local coordinates  $(x^1, \dots, x^n)$  on  $U \subset M$ .

Let  $X = \sum_{i=1}^n X^i \frac{\partial}{\partial x^i}$  and  $Y = \sum_{i=1}^n Y^i \frac{\partial}{\partial x^i}$  be vector fields

The Lie bracket  $[X, Y]$  is the vector field

$$[X, Y] = \sum_{i,j=1}^n \left( X^i \frac{\partial Y^j}{\partial x^i} - Y^i \frac{\partial X^j}{\partial x^i} \right) \frac{\partial}{\partial x^j}$$

These terms "look like" derivatives of vector fields. Are they the same as the corresponding covariant derivatives?

Translating in terms of  $\nabla$ , what we are asking is

$$\otimes \otimes [X, Y] = \nabla_X Y - \nabla_Y X \quad ?$$

Def The function  $T: \mathcal{X}(M) \times \mathcal{X}(M) \rightarrow \mathcal{X}(M)$

$$T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y]$$

is the torsion of  $\nabla$ . The connection  $\nabla$  is torsion free if  $T=0$ , which is the same as saying  $\otimes \otimes$  holds.

Remark: for any connection  $\nabla$  in  $TM$ , the torsion  $T$  is tensorial, i.e.  $C^\infty(M)$ -linear in each factor. Therefore, by Homework 3 Problem 1,  $T$  corresponds to a section of  $(T^*M) \otimes (T^*M) \otimes TM$ , and we can write

$$T = \sum_{i,j,k=1}^n T_{ij}^k dx^i \otimes dx^j \otimes \frac{\partial}{\partial x^k} \quad (\text{in local coordinates})$$

Theorem Let  $(M, g)$  be a Riemannian manifold. There is a unique connection  $\nabla$  in  $TM$  that is compatible with  $g$  and torsion free. This  $\nabla$  is called the **Levi-Civita connection**.

Unfortunately no proof of this theorem seems to be very intuitive.

Proof. We first prove uniqueness by getting a formula for  $\nabla$  in terms of  $g$  and intrinsic operations on  $M$ . Then we check that this formula actually has the right properties.

Suppose that  $\nabla$  is compatible with  $g$  and torsion free. Let  $X, Y, Z$  be vector fields on  $M$ . By compatibility with  $g$ :

$$(I) \quad X \cdot g(Y, Z) = g(\nabla_X Y, Z) + g(Y, \nabla_X Z)$$

$$(II) \quad Y \cdot g(Z, X) = g(\nabla_Y Z, X) + g(Z, \nabla_Y X)$$

$$(III) \quad Z \cdot g(X, Y) = g(\nabla_Z X, Y) + g(X, \nabla_Z Y)$$

Now consider  $(I) + (II) - (III)$

$$X \cdot g(Y, Z) + Y \cdot g(Z, X) - Z \cdot g(X, Y)$$

$$= \left\{ \begin{array}{l} {}^1 g(\nabla_X Y, Z) + {}^2 g(Y, \nabla_X Z) \\ + {}^3 g(\nabla_Y Z, X) + {}^4 g(Z, \nabla_Y X) \\ - {}^5 g(\nabla_Z X, Y) - {}^6 g(X, \nabla_Z Y) \end{array} \right\}$$

1 2 3 4 5 6  
just to keep  
track of terms

$$= \left\{ \begin{array}{l} {}^1 g(\nabla_X Y, Z) + {}^4 g(\nabla_Y X, Z) \\ + {}^2 g(\nabla_X Z, Y) - {}^5 g(\nabla_Z X, Y) \\ + {}^3 g(\nabla_Y Z, X) - {}^6 g(\nabla_Z Y, X) \end{array} \right\}$$

Rearrange terms  
and use that  
 $g$  is symmetric

$$= \left\{ \begin{array}{l} g(\nabla_X Y + \nabla_Y X, Z) \\ + g(\nabla_X Z - \nabla_Z X, Y) \\ + g(\nabla_Y Z - \nabla_Z Y, X) \end{array} \right\}$$

by bilinearity of  $g$

$$= \left\{ \begin{array}{l} g(\nabla_X Y + \nabla_X Y + [Y, X], Z) \\ + g([X, Z], Y) \\ + g([Y, Z], X) \end{array} \right\}$$

By torsion freeness  
e.g.  
 $\nabla_Y X = \nabla_X Y + [Y, X]$

Now there is only one term that involves  $\nabla$ :

$$\begin{aligned} & X \cdot g(Y, Z) + Y \cdot g(Z, X) - Z \cdot g(X, Y) \\ &= 2g(\nabla_X Y, Z) + g([Y, X], Z) + g([X, Z], Y) \\ &\quad + g([Y, Z], X) \end{aligned}$$

This shows that  $g(\nabla_X Y, Z)$  can be solved for in terms of the metric  $g$ , the Lie bracket  $[,]$ , and the action of vector fields on functions. We can now conclude that  $\nabla$  is unique, if it exists. For suppose that  $\nabla^1$  and  $\nabla^2$  are two such connections. Then

$$g(\nabla_X^1 Y, Z) = g(\nabla_X^2 Y, Z) \quad \text{for all } X, Y, Z \in \mathfrak{X}(M)$$

$$\text{Then } g(\nabla_X^2 Y - \nabla_X^1 Y, Z) = 0 \quad \text{for all } X, Y, Z \in \mathfrak{X}(M)$$

Since  $g$  is non-degenerate, the only vector field that is orthogonal to all vector fields  $Z$  is 0. Hence  $\nabla_X^2 Y - \nabla_X^1 Y = 0$ .

For existence, the idea is to **actually define**  $\nabla$  by the relation

$$\begin{aligned} 2g(\nabla_X Y, Z) &= X \cdot g(Y, Z) + Y \cdot g(Z, X) - Z \cdot g(X, Y) \\ &\quad - \left[ g([Y, X], Z) + g([X, Z], Y) + g([Y, Z], X) \right] \end{aligned}$$

Sketch: First, show that, for fixed  $X$  and  $Y$ , the right hand side is tensorial in the  $Z$  variable. Thus the right hand side defines a section of  $T^*M$  for fixed  $X$  and  $Y$ .

Since  $g$  is nondegenerate, any section  $\alpha \in \Gamma(M, T^*M)$  is of the form  $\alpha = g(A, -)$  for some section  $A \in \Gamma(M, T^*M)$

So for fixed  $X$  and  $Y$ , there is a vector field  $A$  such that  $g(A, Z) =$  Right hand side

If we define  $\nabla_X Y = \frac{1}{2}A$  then the defining relation is satisfied. Now we need to check that  $\nabla_X Y$  actually defines a connection:

(1) linear in  $Y$ , (2) Leibniz rule, (3) tensorial in  $X$ .

These follow from direct manipulation of the defining relation.  

Now that we know that the Levi-Civita exists, we can compute it in local coordinates  $(x^1, \dots, x^n)$ . Let  $\frac{\partial}{\partial x^i}$  be the coordinate vector fields. They are also a **frame for  $TM$** , so the connection is determined by the **Christoffel symbols**

$$\Gamma_{ij}^k \text{ such that } \nabla_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} = \sum_{k=1}^n \Gamma_{ij}^k \frac{\partial}{\partial x^k}$$

The defining relation for  $\nabla$  drastically simplifies when  $X, Y, Z$  are coordinate vector fields  $\frac{\partial}{\partial x^i}$ , since the Lie bracket terms vanish:

$$\left[ \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right] = 0$$

$$\begin{aligned} 2g\left(\nabla_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j}, \frac{\partial}{\partial x^k}\right) &= \frac{\partial}{\partial x^i} g\left(\frac{\partial}{\partial x^j}, \frac{\partial}{\partial x^k}\right) + \frac{\partial}{\partial x^j} g\left(\frac{\partial}{\partial x^k}, \frac{\partial}{\partial x^i}\right) \\ &\quad - \frac{\partial}{\partial x^k} g\left(\frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j}\right) \end{aligned}$$

Since  $g\left(\frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j}\right) = g_{ij}$ , we get

$$2g\left(\sum_{\ell} \Gamma_{ij}^{\ell} \frac{\partial}{\partial x^{\ell}}, \frac{\partial}{\partial x^k}\right) = \frac{\partial g_{jk}}{\partial x^i} + \frac{\partial g_{ki}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^k}$$

$$\sum_{\ell} 2g_{\ell k} \Gamma_{ij}^{\ell} = \frac{\partial g_{jk}}{\partial x^i} + \frac{\partial g_{ki}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^k}$$

$$\Gamma_{ij}^{\ell} = \sum_k \frac{1}{2} g^{\ell k} \left( \frac{\partial g_{jk}}{\partial x^i} + \frac{\partial g_{ki}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^k} \right)$$

where  $(g^{\ell k})$  is the inverse matrix of  $(g_{\ell k})$

That is:  $\sum_k g_{ik} g^{jk} = \delta_i^j$

Remark The fact that  $\nabla$  is torsion free boils down to

$$\Gamma_{ij}^k = \Gamma_{ji}^k \quad (\Gamma \text{ is symmetric in } i \text{ and } j)$$